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Boundary Control of Nonlinear ODE/Wave PDE Systems with Spatially-Varying Propagation Speed

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Abstract-We consider the boundary control of a nonlinear ODE actuated through a wave equation whose propagation speed is spatially-varying. The ODE state is driven by the uncontrolled boundary of the wave equation. We design a nonlinear backstepping compensator to enable global asymptotic stability of the closed-loop system. We deduce the controller design and the stability proof by introducing a two-step backstepping transformation. The first transformation recasts the original system into a coupled 2×2 first-order hyperbolic system with spatially-varying coefficients cascading into a nonlinear ODE. The second transformation is used in the design of a compensator for the resulting cascaded system. Our design offers a global stability result that is guaranteed assuming that the spatiallyvarying propagation speed is continuously differentiable and positive. Moreover, for nonlinear systems, our result is the first contribution enabling actual compensation of actuator delays governed by a coupled first-order hyperbolic PDEs induced by a wave PDE dynamics with spatially-varying propagation speed. The validity of the proposed controller is illustrated by the benchmark system controlled via a cable.

Index Terms—Nonlinear system, wave PDE, predictor feedback control, spatially-varying coefficients.

I. INTRODUCTION

Recent studies using backstepping control technique have enabled the stabilization of nonlinear ordinary differential equation (ODE) systems with input delays that depend on the ODE state in [1], [2], [3], [4], as well as the uncontrolled-or controlled-boundary value of the partial differential equation (PDE) state in [5], [6], [7], [8]. Later, the method has been extended to deal with the stabilization problem of nonlinear systems with actuator dynamics governed by wave PDE with moving boundary that depends on the ODE state [9], [10], [11]. The method has also been employed to control transport PDE-ODE cascades with delayed input [12], as well as with state-dependent propagation speed [13], [14].

The result of [15] on the boundary feedback control of PDE-ODE cascaded systems highlighted the potential of the PDE control for various physical systems. Along the same lines, in [16], the stabilization of a linear ODE whose actuator dynamics is governed by a first-order linear hyperbolic PDE

is achieved via backstepping design. Using a two-step backstepping method, boundary control of linear ODE with linear 2×2 hyperbolic systems with spatially-varying coefficients and dynamic boundary conditions has been established in [18] and results dealing with the inverse optimal control for systems with input delays can be found in [19], [20], [21], [22]. The cascaded system consisting of an ODE actuated via a hyperbolic PDE is relevant to many engineering problems including metal rolling processes [23], and metal cutting processes [24], vehicular traffic flow [25], moisture in convective flows [26], transport phenomena in gasoline engines [27], [28], commercial fuels by blending [29].

The particular case of actuator dynamics governed by a wave PDE is quite interesting and has been proven to enable stabilization of stick-slip instabilities and bit-bouncing phenomena in oil drilling processes [13], [30], [31], [32], [33]. These phenomena lead to growing torsional and axial vibrations resulting from the complex interaction between the drill bit and the deeply cracked rock when operating drill strings [34]. In fact, neglecting the damping coefficients between the mud and the pipe, the axial and torsional excitations of the drill string can be described by a wave PDE cascading into a nonlinear ODE governing the dynamic boundary condition at the bit-rock point of contact. Through linear and bilinear matrix inequalities techniques, feedback controllers are established guaranteeing ultimate boundedness of the system trajectories and leading consequently to the suppression of harmful dynamics in drilling system [35]. The dynamics of a flexible cable crane with a load can be expressed by a wave PDE/ nonlinear ODE cascaded system [36]. One should mention recent advances achieved in stabilizing 2×2 coupled hyperbolic PDEs in cascade with linear ODEs to cancel oscillations of tension and cage in dual-cable mining elevators [37] and deep-sea construction [38].

In the present work, we deal with a problem that is similar but not equivalent to the model used to stabilize oscillations in drilling systems. The main goal of this contribution is to design a compensator for the delay induced by the wave PDE with a spatially distributed propagation speed. Though many similarities might be found in the existing literature regarding the stability proof rational, the derived predictor state cannot be obtained by extending any of the existing results based on PDE backstepping design. The key difference between the structure of the considered system and the one investigated in [9] arises after transforming the original wave PDE into a linear 2×2 hyperbolic system. Clearly, in [9], the resulting transport PDEs are decoupled while the presence of a spatially distributed coefficient induces strongly coupled hyperbolic

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PDEs with spatially distributed coefficients cascading with a nonlinear ODE in our case. The strong coupling makes the design of an actuator dynamics compensator non-trivial and different from [9] and [10], [11], which deal with constant propagation speed. Boundary controllers for consisting of a linear ODE in cascade with coupled hyperbolic PDEs can be found in [37], [38]. Mathematically, *a novel two-step backstepping transformation* is employed to derive a non-standard target system whose stability is established using a Lyapunov argument. The resulting boundary controller is a predictor-feedback control law, which compensates the wave actuator dynamics and guarantees global asymptotic stability of the closed-loop system.

This paper is organized as follows: the system's description, the main result are presented in Section II. The transformation of the original system into a 2×2 coupled linear hyperbolic PDEs is in Section III. A first backstepping transformation is introduced in Section IV-A. A second backstepping transformation is employed to design a compensator for the resulting decoupled system in Section IV-B. The stability analysis of the target system is established in Section V. Stability of the original system is stated in Section VI. Finally, an example is provided in Section VII, and concluding remarks are emphasized in Section VIII.

Notation. We use the common definitions of class \mathscr{K} , \mathscr{K}_{∞} , $\mathscr{K}\mathscr{L}$ functions from [29]. For an n-vector, $|\cdot|$ denotes the usual Euclidean norm. For a scalar function $u(\cdot,t)$, we denote with $||u(t)||_{\infty}$ its supremum norm, i.e. $||u(t)||_{\infty} = \sup_{x \in [0,L]} |u(x,t)|$.

II. PROBLEM STATEMENT AND MAIN RESULT

A. System description

We consider a cascaded system consisting of a nonlinear ODE whose actuation path is governed by a wave PDE with spatially-varying propagation speed. The cascaded system has the structure

$$\dot{X}(t) = f(X(t), u(0, t)),$$
 (1)

$$\partial_{tt}u(x,t) = v(x)\partial_{xx}u(x,t), \qquad (2)$$

$$\partial_x u(0,t) = 0, \tag{3}$$

$$\partial_x u(L,t) = U(t),$$
 (4)

where $t \ge 0, 0 \le x \le L$ and $X \in \mathbb{R}^n, u \in \mathbb{R}, U \in \mathbb{R}$ are ODE state, PDE state, and control input, respectively, and $f : \mathbb{R}^n \times \mathbb{R} \to \mathbb{R}^n$ is locally Lipschitz with $f(0,0) = 0, v : [0,L] \to \mathbb{R}_+$.

We design a predictor control that stabilizes the PDE/ODE cascaded system, under the following Assumptions.

Assumption 1: Propagation speed v(x) is continuously differentiable and positive for all $x \in [0, L]$.

Remark 1: Denote

$$\underline{v} = \inf_{x \in [0,L]} v(x), \quad \overline{v} = \sup_{x \in [0,L]} v(x), \tag{5}$$

for all $x \in [0, L]$, the following holds

$$\overline{\upsilon} \ge v(x) \ge \underline{\upsilon} > 0. \tag{6}$$

Assumption 2: System $\dot{X} = f(X, \kappa(X) + \upsilon)$ is input-to-state stable with respect to υ and the function $\kappa : \mathbb{R}^n \to \mathbb{R}$ is continuously differentiable with locally Lipschitz derivative $\frac{\partial \kappa(X)}{\partial X}$ and it satisfies $\kappa(0) = 0$.

Now, consider $v \in \mathbb{R}$ and define the variables

$$Z(t) = \begin{bmatrix} X(t) \\ u(0,t) \end{bmatrix}, \quad \varphi(Z(t),\upsilon) = \begin{bmatrix} f(X(t),u(0,t)) \\ \upsilon \end{bmatrix}.$$
(7)

Assumption 3: System $\dot{Z} = \varphi(Z, \upsilon)$ is strongly forward/backward complete with respect to υ , that is, there exist smooth positive definite functions R_1, R_2 and class \mathscr{K}_{∞} functions $\alpha_1, \dots, \alpha_6$ such that for all $Z \in \mathbb{R}^{n+1}$ and $\upsilon \in \mathbb{R}$,

$$\alpha_1(|Z|) \le R_1(Z) \le \alpha_2(|Z|) \tag{8}$$

$$\frac{\partial R_1(Z)}{\partial Z}\varphi(Z,\upsilon) \le R_1(Z) + \alpha_3(|\upsilon|) \tag{9}$$

$$\alpha_4(|Z|) \le R_2(Z) \le \alpha_5(|Z|) \tag{10}$$

$$-\frac{\partial R_2(Z)}{\partial Z}\varphi(Z,\upsilon) \le R_2(Z) + \alpha_6(|\upsilon|).$$
(11)

Assumption 4: System $\dot{Z} = \varphi(Z, \mu(Z) + \upsilon)$ is strongly backward complete with respect to υ , that is, there exist a smooth positive definite function R_3 and class \mathscr{K}_{∞} functions α_7 , α_8 , α_9 such that for $Z \in \mathbb{R}^{n+1}$ and $\upsilon \in \mathbb{R}$,

$$\alpha_7(|Z|) \leq R_3(Z) \leq \alpha_8(|Z|) \quad (12)$$

$$-\frac{\partial R_3(Z)}{\partial Z}\varphi(Z,\mu(Z)+\upsilon) \leq R_3(Z)+\alpha_9(|\upsilon|).$$
(13)

Remark 2: The wave PDE described by (2)–(4) induces an input delay on the nonlinear dynamics of (1). For instance, an actuator dynamics governed by a pure transport PDE with spatially varying coefficient, defined as

$$\dot{X}(t) = f(X(t), u(0, t)),$$
(14)

$$\partial_t u(x,t) = v(x)\partial_x u(x,t), \quad u(L,t) = U(t), \tag{15}$$

is equivalent to the nonlinear system with input delay $\dot{X}(t) = f(X(t), U(\phi(t)))$, where $\phi(t) = t - \int_0^L v^{-1}(s) ds$.

B. Main results

Suppose the existence of a nominal controller $\kappa : \mathbb{R}^n \to \mathbb{R}$ that stabilizes the delay-free plant, namely, the control law κ is such that the closed-loop system $\dot{X}(t) = f(X(t), \kappa(X(t)))$ is globally asymptotically stable. The predictor feedback control for system (1)–(4) is given by

$$\begin{split} U(t) &= -c_1 \frac{e^{\int_0^L \frac{v'(r)}{4v(r)} dr}}{2\sqrt{v(L)}} (p_2(L,t) - \kappa(p_1(L,t))) \\ &+ \frac{e^{\int_0^L \frac{v'(r)}{4v(r)} dr}}{2\sqrt{v(L)}} \frac{\partial \kappa(p_1(L,t))}{\partial p_1} f(p_1(L,t), p_2(L,t)) \\ &- \frac{1}{2\sqrt{v(L)}} \left(\partial_t u(L,t) - \sqrt{v(L)} \partial_x u(L,t) \right) \\ &+ \frac{1}{2\sqrt{v(L)}} \int_0^L K_{11}(L,s) \left(\partial_t u(s,t) + \sqrt{v(s)} \partial_s u(s,t) \right) ds \\ &+ \frac{1}{2\sqrt{v(L)}} \int_0^L K_{12}(L,s) \left(\partial_t u(s,t) - \sqrt{v(s)} \partial_s u(s,t) \right) ds, \end{split}$$
(16)

where $p_1 \in \mathbb{R}^n, p_2 \in \mathbb{R}$ are given by

$$p_1(x,t) = X(t) + \int_0^x \frac{f(p_1(y,t), p_2(y,t))}{\sqrt{v(y)}} dy, \qquad (17)$$

$$p_{2}(x,t) = u(0,t) + \int_{0}^{x} \left(\frac{1}{\sqrt{v(y)}} - \int_{y}^{x} \frac{K_{11}(\sigma, y)}{\sqrt{v(\sigma)}} d\sigma \right)$$
$$\times e^{-\int_{0}^{y} \frac{v'(r)}{4v(r)} dr} (\partial_{t}u(y,t) + \sqrt{v(y)} \partial_{y}u(y,t)) dy$$
$$- \int_{0}^{x} \int_{0}^{x} \frac{K_{12}(\sigma, y)}{\sqrt{\sigma}} d\sigma$$

$$\sum_{v \in T_{0}}^{y} \sqrt{v(\sigma)} \sqrt{v(\sigma)}$$

$$\times e^{-\int_{0}^{y} \frac{v'(r)}{4v(r)} dr} (\partial_{t} u(y,t) - \sqrt{v(y)} \partial_{y} u(y,t)) dy,$$
(18)

for all $x \in [0, L]$ with initial conditions as

$$p_{1}(x,0) = X(0) + \int_{0}^{x} \frac{f(p_{1}(y,0), p_{2}(y,0))}{\sqrt{v(y)}} dy$$
(19)
$$p_{2}(x,0) = u(0,0) + \int_{0}^{x} \left(\frac{1}{\sqrt{v(y)}} - \int_{y}^{x} \frac{K_{11}(\sigma,y)}{\sqrt{v(\sigma)}} d\sigma\right)$$
$$\times e^{-\int_{0}^{y} \frac{y'(r)}{4v(r)} dr} (\partial_{t}u(y,0) + \sqrt{v(y)} \partial_{y}u(y,0)) dy$$
$$- \int_{0}^{x} \int_{y}^{x} \frac{K_{12}(\sigma,y)}{\sqrt{v(\sigma)}} d\sigma$$
$$\times e^{-\int_{0}^{y} \frac{y'(r)}{4v(r)} dr} (\partial_{t}u(y,0) - \sqrt{v(y)} \partial_{y}u(y,0)) dy.$$
(20)

The gain c_1 in (16) is an arbitrary constant while the kernel gains K_{11} and K_{12} are solutions to the following gain kernel PDEs:

$$\mathscr{A}(x)\partial_x K(x,s) + \partial_s (K(x,s)\mathscr{A}(s)) = K(x,s)\mathscr{B}(s)$$
(21)

$$K(x,x)\mathscr{A}(x) - \mathscr{A}(x)K(x,x) = \mathscr{B}(x)$$
(22)

$$K_{11}(x,0) = K_{12}(x,0) \tag{23}$$

$$K_{21}(x,0) = K_{22}(x,0) \tag{24}$$

where (21) is defined on $\{(x,s): 0 \le s \le x \le 1\}$, and

$$K(x,s) = \begin{bmatrix} K_{11}(x,s) & K_{12}(x,s) \\ K_{21}(x,s) & K_{22}(x,s) \end{bmatrix} \in \mathbb{R}^{2 \times 2},$$
(25)

and

$$\mathscr{A}(x) = \begin{bmatrix} \sqrt{v(x)} & 0\\ 0 & -\sqrt{v(x)} \end{bmatrix}, \mathscr{B}(x) = \begin{bmatrix} 0 & -\frac{v'(x)}{4\sqrt{v(x)}}\\ \frac{v'(x)}{4\sqrt{v(x)}} & 0 \end{bmatrix}.$$
(26)

Theorem 1: Consider system (1)–(4) together with the control law (16)–(18), for any initial condition $u(\cdot,0) \in C_1[0,1]$, $u_t(\cdot,0) \in C[0,1]$, which is compatible with the feedback law (16)–(18) and which is such that $u_x(0,0) = 0$. Under Assumptions 1–4, the closed-loop system has a unique solution $X(t) \in C_1([0,\infty), \mathbb{R}^n)$, $(u(\cdot,t), u_t(\cdot,t)) \in C([0,\infty), C_1[0,1] \times C[0,1])$. Moreover, there is a \mathcal{KL} function $\overline{\beta}$ such that

$$\Omega(t) \leq \beta(\Omega(0), t) \tag{27}$$

$$\Omega(t) = |X(t)| + ||u(t)||_{\infty} + ||u_t(t)||_{\infty} + ||u_x(t)||_{\infty}, \quad (28)$$

for all $t \ge 0$.

Stability analysis of the closed-loop system will be derived in the following five steps:

- 1) Introduction of two transformations to recast the system into a coupled 2×2 hyperbolic PDEs with spatially-varying coefficients is shown in Section III.
- 2) Removal of the internal PDE states coupling acting on the resulting 2×2 the hyperbolic PDEs via a first backstepping transformation is shown in Section IV-A.
- The predictor feedback of the equivalent decoupled PDE/nonlinear ODE target system obtained from a second backstepping transformation is introduced in Section IV-B.
- 4) Stability analysis for the target system is provided in Section V.
- 5) Proof of stability of the original cascaded system is established in Section VI.

Remark 3: The proof of the well-posedness of the gain kernel PDEs defined in (21)–(24) can be found in [39] which deals with the stabilization of a 2×2 coupled hyperbolic PDE using backstepping technique.

III. FROM A WAVE PDE/ODE CASCADED SYSTEM TO A COUPLED 2×2 hyperbolic PDE system with spatially-varying coefficients

In this section we employ two changes of coordinates denoted *Transformation I* and *Transformation II* in order to map system (1)–(4) into a suitable coupled first-order hyperbolic system cascading into a nonlinear ODE.

A. Coupled hyperbolic system: Transformation I

First, we introduce the following change of coordinate

$$\overline{\zeta}(x,t) = \partial_t u(x,t) + \sqrt{v(x)} \,\partial_x u(x,t), \tag{29}$$

$$\overline{\eta}(x,t) = \partial_t u(x,t) - \sqrt{v(x)} \,\partial_x u(x,t) \,, \tag{30}$$

which in reverse is written as

$$\partial_t u(x,t) = \frac{\overline{\zeta}(x,t) + \overline{\eta}(x,t)}{2},$$
(31)

$$\partial_x u(x,t) = \frac{\overline{\zeta}(x,t) - \overline{\eta}(x,t)}{2\sqrt{\nu(x)}}.$$
(32)

Taking the time and spatial derivatives of (29) and (30), we map the original system (2)–(4) into the following coupled hyperbolic PDE cascading with the nonlinear ODE defined in (1)(see Fig. 1).

$$\dot{X} = f(X, u(0, t)) \tag{33}$$

$$\partial_t \overline{\xi}(x,t) = \mathscr{A}(x)\partial_x \overline{\xi}(x,t) + \mathscr{B}_0(x)\overline{\xi}(x,t)$$
(34)

$$\partial_t u(0,t) = \zeta(0,t) \tag{35}$$

$$\overline{\eta}(0,t) = \zeta(0,t) \tag{36}$$

$$\zeta(L,t) = \overline{\eta}(L,t) + 2\sqrt{\nu(L)U(t)}, \qquad (37)$$



Fig. 1: Equivalent 2×2 coupled hyperbolic PDEs system cascading into a nonlinear ODE

where

$$\overline{\xi}(x,t) = \begin{bmatrix} \overline{\zeta}(x,t) \\ \overline{\eta}(x,t) \end{bmatrix}, \mathscr{B}_0(x) = \begin{bmatrix} \frac{-\nu'(x)}{4\sqrt{\nu(x)}} & \frac{-\nu'(x)}{4\sqrt{\nu(x)}} \\ \frac{\nu'(x)}{4\sqrt{\nu(x)}} & \frac{\nu'(x)}{4\sqrt{\nu(x)}} \end{bmatrix}, \quad (38)$$

and $\mathscr{A}(x)$ is given by (26).

B. Coupled hyperbolic system: Transformation II

Applying the state transformations

$$\zeta(x,t) = e^{-\int_0^x \frac{v'(r)}{4v(r)} dr} \overline{\zeta}(x,t),$$
(39)

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$$\eta(x,t) = e^{-\int_0^x \frac{v'(r)}{4v(r)} dr} \overline{\eta}(x,t),$$
(40)

system (33)-(37) is rewritten in the following form

$$\dot{Z}(t) = \varphi(Z(t), \zeta(0, t)) \tag{41}$$

$$\partial_t \xi(x,t) = \mathscr{A}(x)\partial_x \xi(x,t) + \mathscr{B}(x)\xi(x,t)$$
(42)

$$\zeta(0,t) = \eta(0,t) \tag{43}$$

$$\zeta(L,t) = \eta(L,t) + 2e^{-\int_0^L \frac{v'(r)}{4v(r)}dr} \sqrt{v(L)} U(t), \qquad (44)$$

where $\xi(x,t) = \begin{bmatrix} \zeta(x,t), & \eta(x,t) \end{bmatrix}^T$, and $\mathscr{A}(x)$ and $\mathscr{B}(x)$ are given by (26).

IV. BACKSTEPPING TRANSFORMATIONS

A. First step backstepping transformation

In this section, we employ a first backstepping transformation to system (41)–(44) in order to remove the internal PDE state coupling terms of the two propagating waves as shown in Fig. 3:

$$\boldsymbol{\omega}(x,t) \quad = \quad \boldsymbol{\xi}(x,t) - \int_0^x K(x,s)\boldsymbol{\xi}(s,t)ds, \tag{45}$$

for all $0 \le x \le L$, $t \ge 0$. The gain kernel matrix K(x,s) is a solution to the kernel equations (21)–(24), where (21) is defined on $\{(x,s): 0 \le s \le x \le L\}$. Following [39], it can be verified that (45) has a bounded inverse defined as

$$\xi(x,t) = \omega(x,t) + \int_0^x L(x,s)\omega(s,t)ds, \qquad (46)$$

for all $0 \le x \le L, t \ge 0$. Moreover, the inverse gain kernel matrix

$$L(x,s) = \begin{bmatrix} L_{11}(x,s) & L_{12}(x,s) \\ L_{21}(x,s) & L_{22}(x,s) \end{bmatrix} \in \mathbb{R}^{2 \times 2},$$
(47)



Fig. 2: First-step backstepping transformation removes the internal coupling terms

is the solution of the well-posed kernel equations written below

$$\partial_s L(x,s) \mathscr{A}(x)) - \mathscr{A}(x) \partial_x L(x,s) = -\mathscr{B}(x) L(x,s)$$
 (48)

$$L(x,x)\mathscr{A}(x) - \mathscr{A}(x)L(x,x) = \mathscr{B}(x)$$
(49)

$$L_{11}(x,0) = L_{12}(x,0)$$
(50)

$$L_{21}(x,0) = L_{22}(x,0),$$
 (51)

where $\mathscr{A}(x)$ and $\mathscr{B}(x)$ are defined by (26).

Differentiating (45) with respect to time t and space x, it can be straightforwardly established that system (41)–(44) maps into the following decoupled PDE/ODE cascaded system

$$\dot{Z}(t) = \varphi(Z(t), \omega_1(0, t))$$
(52)

$$\partial_t \omega(x,t) = \mathscr{A}(x) \partial_x \omega(x,t) \tag{53}$$

$$\boldsymbol{\omega}_2(0,t) = \boldsymbol{\omega}_1(0,t) \tag{54}$$

$$\omega_{1}(L,t) = \eta(L,t) + 2e^{-\int_{0}^{L} \frac{v'(r)}{4v(r)}dr} \sqrt{v(L)} U(t)$$

$$-\int_{0}^{L} K_{11}(L,s)\zeta(s,t)ds$$

$$-\int_{0}^{L} K_{12}(L,s)\eta(s,t)ds, \qquad (55)$$

for $0 \le x \le L, t \ge 0$, and $\omega(x,t) = [\omega_1(x,t), \omega_2(x,t)]^T$ if the gain kernel matrix satisfies (21)–(24). The ODE dynamics is driven by φ defined in (7).

B. Second-step backstepping transformation

From a nominal controller $\kappa : \mathbb{R}^n \to \mathbb{R}$ that globally asymptotically stabilizes the original delay-free nonlinear ODE, we define the following function

$$\mu(\boldsymbol{\chi}) = -c_1(\boldsymbol{\chi}_2 - \boldsymbol{\kappa}(\boldsymbol{\chi}_1)) + \frac{\partial \boldsymbol{\kappa}(\boldsymbol{\chi}_1)}{\partial \boldsymbol{\chi}_1} f(\boldsymbol{\chi}_1, \boldsymbol{\chi}_2), \quad (56)$$

where $c_1 > 0$ is an arbitrary gain constant and $\chi = [\chi_1, \chi_2] \in \mathbb{R}^n \times \mathbb{R}$. Denote the vector functions p(x,t) and q(x,t) as

$$p(x,t) = Z(t) + \int_0^x \frac{\varphi(p(y,t), \omega_1(y,t))}{\sqrt{v(y)}} dy,$$
 (57)

where $p(x,t) = [p_1(x,t), p_2(x,t)]^T$, with the initial condition

$$p(x,0) = Z(0) + \int_0^x \frac{\varphi(p(y,0), \omega_1(y,0))}{\sqrt{\nu(0)}} dy, \qquad (58)$$

$$\frac{-\sqrt{v(x)}}{\left(\frac{\sqrt{v(x)}}{L}\right)} \xrightarrow{\lambda(x,t)} \overline{\omega(L,t) = 0}$$

Fig. 3: Second-step backsteping transformation removes the boundary coupling at x = L, and designs a compensator for the resulting cascaded system

and

$$q(x,t) = Z(t) - \int_0^x \frac{\varphi(q(y,t), \omega_2(y,t))}{\sqrt{v(y)}} dy,$$
 (59)

where $q(x,t) = [q_1(x,t), q_2(x,t)]^T$, with the initial condition

$$q(x,0) = Z(0) - \int_0^x \frac{\varphi(q(y,0), \omega_2(y,0))}{\sqrt{v(0)}} dy.$$
 (60)

Lemma 1 (Second-Step Backstepping Transform): The following backstepping transformations

$$\boldsymbol{\varpi}(\boldsymbol{x},t) = \boldsymbol{\omega}_{1}(\boldsymbol{x},t) - \boldsymbol{\mu}(\boldsymbol{p}(\boldsymbol{x},t)), \quad (61)$$

$$\lambda(x,t) = \omega_2(x,t) - \mu(q(x,t)), \qquad (62)$$

where μ is defined in (56), and p(x,t), q(x,t) are given as (57), (59), respectively, and U(t) is

$$U(t) = \frac{e^{\int_0^L \frac{y'(s)}{4v(s)}ds}}{2\sqrt{v(L)}} \bigg(\mu(p(L,t)) - \eta(L,t) + \int_0^L \big(K_{11}(L,s)\zeta(s,t) + K_{12}(L,s)\eta(s,t)\big)ds\bigg), \quad (63)$$

map system (52)-(55) into the target system

$$\dot{Z}(t) = \varphi\left(Z(t), \varpi(0, t) + \mu(Z(t))\right)$$
(64)

$$\partial_t \boldsymbol{\varpi}(x,t) = \sqrt{v(x)} \,\partial_x \boldsymbol{\varpi}(x,t)$$
 (65)

$$\partial_t \lambda(x,t) = -\sqrt{\nu(x)} \,\partial_x \lambda(x,t)$$
 (66)

$$\lambda(0,t) = \boldsymbol{\varpi}(0,t) \tag{67}$$

$$\boldsymbol{\varpi}(L,t) = 0. \tag{68}$$

The schematic of the resulting target system is depicted in Fig. 3.

Proof. It is easy to obtain (64), (67), (68). We will prove relation (65). Differentiating (57) with respect to t and x, we get

$$\partial_t p(x,t) = \varphi(Z(t), \omega_1(0,t)) + \int_0^x \frac{\partial_p \varphi(p(y,t), \omega_1(y,t)) \partial_t p(y,t)}{\sqrt{v(y)}} dy + \int_0^x \frac{\partial_{\omega_1} \varphi(p(y,t), \omega_1(y,t)) \partial_t \omega_1(y,t)}{\sqrt{v(y)}} dy, \qquad (69)$$

and

$$\sqrt{v(x)}\partial_x p(x,t) = \int_0^x \partial_p \varphi(p(y,t), \omega_1(p(y,t))) \partial_y p(y,t) dy + \int_0^x \partial_{\omega_1} \varphi(p(y,t), \omega_1(p(y,t))) \partial_y \omega_1(y,t) dy + \varphi(Z(t), \omega_1(0,t)),$$
(70)

respectively. Defining $H(x,t) = \partial_t p(x,t) - \sqrt{v(x)} \partial_x p(x,t)$, and combining (69) and (70), and by (53), we arrive at

$$H(x,t) = \int_0^x \frac{\partial_p \varphi(p(y,t), \omega_1(y,t)) H(y,t)}{\sqrt{\nu(y)}} dy.$$
(71)

Differentiating (71) with respect to x, we have

$$\partial_x H(x,t) = \frac{\partial_p \varphi(p(x,t), \omega_1(x,t)) H(x,t)}{\sqrt{v(x)}},$$
(72)

and H(0,t) = 0, which implies that H(x,t) = 0, for all $x \in [0,L]$. Hence, it is clear that

$$\partial_t p(x,t) = \sqrt{v(x)} \partial_x p(x,t).$$
 (73)

Taking the time and the spatial derivatives of (61), and from (53) and (73), we obtain (65). Relation (66) can be deduced similarly, which completes the proof.

Remark 4: Using (39), it is easy to show that (63) can be expressed as follows:

$$U(t) = \frac{e^{\int_0^L \frac{v'(t)}{4v(t)}ds}}{2\sqrt{v(L)}} \mu(p(L,t)) - \frac{1}{2\sqrt{v(L)}}\bar{\eta}(L,t) + \frac{1}{2\sqrt{v(L)}} \int_0^L K_{11}(L,s)\bar{\zeta}(s,t))ds + \frac{1}{2\sqrt{v(L)}} \int_0^L K_{12}(L,s)\bar{\eta}(s,t)ds,$$
(74)

Substituting (29), (30), and (56) into (74) gives the equivalent control action U(t) defined in (16). In addition, it can be deduced that p(x,t) in (57) is equal to $[p_1(x,t), p_2(x,t)]^T$, where $p_1(x,t), p_2(x,t)$ are given by (17), (18), respectively.

Define the vector functions $\pi(x,t)$ and $\iota(x,t)$ as

1.

$$\pi(x,t) = Z(t) + \int_0^x \frac{\varphi(\pi(y,t), \varpi(y,t) + \mu(\pi(y,t)))}{\sqrt{\nu(y)}} dy, \quad (75)$$

where $\pi(x,t) = [\pi_1(x,t), \pi_2(x,t)]^T$, with the initial condition

$$\pi(x,0) = Z(0) + \int_0^x \frac{\varphi(\pi(y,0), \varpi(y,0) + \mu(\pi(y,0)))}{\sqrt{\nu(y)}} dy,$$
(76)

and

$$\iota(x,t) = Z(t) - \int_0^x \frac{\varphi(\iota(y,t), \lambda(y,t) + \mu(\iota(y,t)))}{\sqrt{\nu(y)}} dy, \quad (77)$$

where $\iota(x,t) = [\iota_1(x,t), \iota_2(x,t)]^T$ with the initial condition

$$\iota(x,0) = Z(0) - \int_0^x \frac{\varphi(\iota(y,0), \lambda(y,0) + \mu(\iota(y,0)))}{\sqrt{\nu(y)}} dy, \quad (78)$$

where $\boldsymbol{\varpi}, \lambda, \mu$ are defined in (61), (62), (56), respectively.

$$\omega_1(x,t) = \boldsymbol{\varpi}(x,t) + \boldsymbol{\mu}(\boldsymbol{\pi}(x,t)), \tag{79}$$

$$\omega_2(x,t) = \lambda(x,t) + \mu(\iota(x,t)), \tag{80}$$

where $\pi(x,t)$, $\iota(x,t)$, $0 \le x \le L$, $t \ge 0$, are given as (75), (77), respectively.

The inverse backstepping transformations (79), (80), and the control law (63) transform the target system (64)–(68) into system (52)–(55) and the proof can be derived from straightforward computations.

V. STABILITY ANALYSIS OF THE TARGET SYSTEM

Lemma 2 (Stability of the Target System): Consider system (64)–(68), under Assumptions 1 and 2, there exists a class \mathscr{KL} function β , such that

$$|Z(t)| + \|\boldsymbol{\varpi}(t)\|_{\infty} + \|\boldsymbol{\lambda}(t)\|_{\infty} \leq \beta(|Z(0)| + \|\boldsymbol{\varpi}(0)\|_{\infty} + \|\boldsymbol{\lambda}(0)\|_{\infty}, t),$$
(81)

for all $t \ge 0$.

Proof. We introduce a new variable $z(x,t), x \in [-L,L]$ such that

$$z(x,t) = \begin{cases} \boldsymbol{\varpi}(x,t), & \text{for all } x \in [0,L], \\ \boldsymbol{\lambda}(-x,t), & \text{for all } x \in [-L,0]. \end{cases}$$
(82)

Let $\Gamma_{g,n}(t)$ denote the following norm

$$\Gamma_{g,n}(t) = \int_{-L}^{L} e^{2ng(L+x)} z(x,t)^{2n} dx,$$
(83)

where g > 0 is determined later and *n* is any positive integer. Using integration by parts, the derivative of $\Gamma_{g,n}(t)$ satisfies

$$\begin{split} \dot{\Gamma}_{g,n}(t) &\leq -\int_{-L}^{0} 2n \left(g \sqrt{v(-x)} - \frac{v'(-x)}{4n\sqrt{v(-x)}} \right) \\ &\times e^{2ng(L+x)} z(x,t)^{2n} dx \\ &- \int_{0}^{L} 2n \left(g \sqrt{v(x)} + \frac{v'(x)}{4n\sqrt{v(x)}} \right) e^{2ng(L+x)} z(x,t)^{2n} dx. \end{split}$$
(84)

Under Assumption 2, v(x) is continuously differentiable and positive for all $x \in [0, L]$, denote

$$\pi_0 = \sup_{x \in [0,L]} \frac{\nu'(x)}{\sqrt{\nu(x)}}, \quad \pi_1 = \inf_{x \in [0,L]} \frac{\nu'(x)}{\sqrt{\nu(x)}}.$$
 (85)

Note $n \ge 1$, we have

$$g\sqrt{\nu(-x)} - \frac{\nu'(-x)}{4n\sqrt{\nu(-x)}} \ge \min\left\{g\sqrt{\underline{\nu}} - \frac{\pi_0}{4}, g\sqrt{\underline{\nu}}\right\}$$
(86)

for all $x \in [-L, 0]$, and

$$g\sqrt{v(x)} + \frac{v'(x)}{4n\sqrt{v(x)}} \ge \min\left\{g\sqrt{\underline{v}} + \frac{\pi_1}{4}, g\sqrt{\underline{v}}\right\}, \quad (87)$$

for all $x \in [0, L]$, and \underline{v} is given by (5). Choose

$$g > \max\left\{\frac{\pi_0}{4\sqrt{\underline{\upsilon}}}, \, \frac{-\pi_1}{4\sqrt{\underline{\upsilon}}}\right\},\tag{88}$$

we get

$$\dot{\Gamma}_{g,n}(t) \le -2n\pi_2\Gamma_{g,n}(t), \quad \text{for} \quad t \ge 0,$$
(89)

with $\pi_2 = \min\{\min\{g\sqrt{\underline{\upsilon}} - \frac{\pi_0}{4}, g\sqrt{\underline{\upsilon}}\}, \min\{g\sqrt{\underline{\upsilon}} + \frac{\pi_1}{4}, g\sqrt{\underline{\upsilon}}\}\}$. Using Assumption 2, from (89), it is easy to deduce that (81) holds.

VI. STABILITY OF THE ORIGINAL WAVE PDE/NONLINEAR ODE CASCADED SYSTEM

To establish stability proof of the closed-loop system (1)–(4), (16)–(18), we show the boundedness of predictors, first. Proofs of Lemmas 3–8 are established using systematic developments already presented in [10] and [13].

Lemma 3 (Bound on Forward Predictor): Under Assumptions 1 and 3, there exists a class \mathscr{K}_{∞} function ρ_1 such that the following holds:

$$\sup_{0 \le x \le L} |p(x,t)| \le \rho_1(|Z(t)| + \|\omega_1(t)\|_{\infty}).$$
(90)

Lemma 4 (Bound on Backward Predictor): Under Assumptions 1 and 4, there exists a class \mathcal{K}_{∞} function ρ_2 such that the following holds:

$$\sup_{0 \le x \le L} |q(x,t)| \le \rho_2(|Z(t)| + \|\omega_2(t)\|_{\infty}).$$
(91)

Lemma 5 (Bound on Extended Forward State Predictor): Under Assumptions 1 and 2, there exists a class \mathscr{K}_{∞} function ρ_3 such that the following holds:

$$\sup_{0 \le x \le L} |\pi(x,t)| \le \rho_3(|Z(t)| + \|\varpi(t)\|_{\infty}).$$
(92)

Lemma 6 (Bound on Extended Backward State Predictor): Under Assumptions 1 and 4, there exists a class \mathscr{K}_{∞} function ρ_4 such that the following holds:

$$\sup_{0 \le x \le L} |\iota(x,t)| \le \rho_4(|Z(t)| + \|\lambda(t)\|_{\infty}).$$
(93)

Next, we show equivalence of norms of original and target PDE states in Lemma 7 and Lemma 8.

Lemma 7: (Original PDE State Bounded by Target PDE State) Under Assumptions 2 and 4, consider system (64)–(68), and output maps are (79), (80), then there exists a class \mathcal{K}_{∞} function γ_2 such that the following holds:

$$|Z(t)| + \|\omega_{1}(t)\|_{\infty} + \|\omega_{2}(t)\|_{\infty} \leq \gamma_{2}(|Z(t)| + \|\overline{\omega}(t)\|_{\infty} + \|\lambda(t)\|_{\infty}).$$
(94)

Lemma 8: (Target PDE State Bounded by Original PDE State) Under Assumptions 2 and 4, consider system (52)–(55), and output maps are (61), (62), then there exists a class \mathcal{K}_{∞} function γ_3 such that the following holds:

$$|Z(t)| + \|\boldsymbol{\sigma}(t)\|_{\infty} + \|\boldsymbol{\lambda}(t)\|_{\infty} \le \gamma_{3}(|Z(t)| + \|\boldsymbol{\omega}_{1}(t)\|_{\infty} + \|\boldsymbol{\omega}_{2}(t)\|_{\infty}).$$
(95)

Proof of Theorem 1. Using Lemmas 2, 7 and 8, with the help of (45), (46), we get

$$\begin{aligned} |Z(t)| + \|\xi(t)\|_{\infty} \\ &\leq |Z(t)| + (1+\bar{L})\|\omega(t)\|_{\infty} \\ &\leq (1+\bar{L})\gamma_{2}(|Z(t)| + \|\boldsymbol{\sigma}(t)\|_{\infty} + \|\lambda(t)\|_{\infty}) \\ &\leq (1+\bar{L})\gamma_{2}(\beta(|Z(0)| + \|\boldsymbol{\sigma}(0)\|_{\infty} + \|\lambda(0)\|_{\infty}, t)) \\ &\leq (1+\bar{L})\gamma_{2}(\beta(\gamma_{3}(|Z(0)| + \|\boldsymbol{\omega}_{1}(0)\|_{\infty} + \|\boldsymbol{\omega}_{2}(0)\|_{\infty}), t)) \\ &\leq (1+\bar{L})\gamma_{2}(\beta(\gamma_{3}(\sqrt{2}(|Z(0)| + \|\boldsymbol{\omega}(0)\|_{\infty}), t)) \\ &\leq (1+\bar{L})\gamma_{2}(\beta(\gamma_{3}(\sqrt{2}(1+\bar{K})(|Z(0)| + \|\xi(0)\|_{\infty}), t)), \end{aligned}$$
(96)

where $\overline{L} = \max_{\substack{(x,y)\in[0,L]\times[0,L]\\ \text{Using (7), (29)-(32), (39), (40), and (96), we obtain the following estimate}} |K(x,y)|$.

$$\begin{split} |X(t)| + |u(0,t)| + \|\partial_{t}u(t)\|_{\infty} + \|\partial_{x}u(t)\|_{\infty} \\ &\leq \sqrt{2}|Z(t)| + \frac{\sqrt{2}}{2}(1 + \frac{1}{\sqrt{\upsilon}})\sqrt[4]{\frac{\upsilon}{\upsilon}}\|\xi(t)\|_{\infty} \\ &\leq \max\{\sqrt{2}, \frac{\sqrt{2}}{2}(1 + \frac{1}{\sqrt{\upsilon}})\sqrt[4]{\frac{\upsilon}{\upsilon}}\}(|Z(t)| + \|\xi(t)\|_{\infty}) \\ &\leq (1 + \overline{L})\max\{\sqrt{2}, \frac{\sqrt{2}}{2}(1 + \frac{1}{\sqrt{\upsilon}})\sqrt[4]{\frac{\upsilon}{\upsilon}}\} \\ &\times \gamma_{2}(\beta(\gamma_{3}(\sqrt{2}(1 + \overline{K})(|Z(0)| + \|\xi(0)\|_{\infty}), t))) \\ &\leq (1 + \overline{L})\max\{\sqrt{2}, \frac{\sqrt{2}}{2}(1 + \frac{1}{\sqrt{\upsilon}})\sqrt[4]{\frac{\upsilon}{\upsilon}}\} \\ &\times \gamma_{2}(\beta(\gamma_{3}(\sqrt{2}(1 + \overline{K})(|X(0)| + \|u(0)\|_{\infty} + 2\sqrt[4]{\frac{\upsilon}{\upsilon}}(1 + \sqrt{\overline{\upsilon}})) \\ &\times (\|\partial_{t}u(0)\|_{\infty} + \|\partial_{x}u(0)\|_{\infty})), t)). \end{split}$$

Finally defining a class \mathscr{K}_{∞} function $\overline{\beta}(s,t) = (1 + \overline{L})\max\{\sqrt{2}, \frac{\sqrt{2}}{2}(1 + \frac{1}{\sqrt{\underline{v}}})\sqrt[4]{\frac{\overline{v}}{\underline{v}}}\}\gamma_2(\beta(\gamma_3(\sqrt{2}(1 + \overline{K}) + 2\sqrt[4]{\frac{\overline{v}}{\underline{v}}}(1 + \sqrt{\overline{v}})s),t)))$, we get (27).

Following [11], it can be proved that under Assumptions 1– 4 and $u(\cdot,0) \in C_1[0,1]$, $u_t(\cdot,0) \in C[0,1]$, which is compatible with the feedback law (16)–(18), the closed-loop system has a unique solution $X(t) \in C_1([0,\infty), \mathbb{R}^n)$, $(u(\cdot,t), u_t(\cdot,t)) \in$ $C([0,\infty), C_1[0,L] \times C[0,L])$.

VII. SIMULATION RESULTS

We consider a cable that is made of dynamic material, and a load that is the benchmark system modeled as a point mass M is attached to the lower end of the cable. Small deflections of the cable denoted u(x,t) obeys the planar motion governed by (2)–(4), where $\partial_x u(1,t)$ of the cable at its upper end serves as a control input, and $v(x) = \sqrt{g\rho(1+\sin(x))+Mg}$ is the force in the cable, and $X(t) = [X_1, X_2, X_3]^T$ is the state of the load which satisfies

$$\dot{X}_1 = X_2 + X_3^2 \tag{98}$$

$$\dot{X}_2 = X_3 \tag{99}$$

$$\dot{X}_3 = -X_2 - 2X_3 + u(0,t). \tag{100}$$

Here, ρ is the material's density and g is the gravity constant. Following [40], a nominal design for system (98)–(100) is given as

$$\kappa(X) = -X_3 - (X_1 + 2X_2 + X_3 + 0.25X_2^2)$$
(101)
+ 0.25X_3^2)(1 + 0.5X_3).

The control law for system (2)–(4) cascading with (98)–(100) is given by (16)–(18). A simulation study is performed



Fig. 4: Response of $(X_1(t), X_2(t))$ under the proposed control (solid red line) and with uncompensated control (101) (blue line), state prediction $(p_{11}(L,t), p_{12}(L,t))$ (dashdot red line)



Fig. 5: Dynamics of $(X_3(t), U(t))$ under the proposed control (solid red line) and with uncompensated control (101) (blue line), state prediction $p_{13}(L,t)$ (dashdot red line)

with L = 1m, M = 1kg, $\rho = 2kg/m$, and $g = 9.81m/s^2$, for the initial values of the ODE states given as $X_1(0) =$ $0.5, X_2(0) = -0.6, X_3(0) = 0.7$ and $u_x(x, 0) = 0, u_t(x, 0) = 0.5$ for $x \in [0,1]$ and the gain parameter $c_1 = 1$. Responses of the states (X_1, X_2, X_3) together with the predictor states $(p_{11}(L,t), p_{12}(L,t), p_{13}(L,t))$ under the proposed control law and the nominal control action without wave actuator dynamics compensation are shown in Fig. 4-Fig. 5. It is clear that the designed predictor enables to compute the future values of the real states and the designed predictor-feedback controller stabilizes the system at the setpoint. However, the uncompensated control action (101) cannot achieve stabilization of the state X_1 to the desired setpoint and leads to a bounded dynamics as depicted in Fig.4. The actuator dynamics for the compensated case is depicted in Fig.6, which confirms the pertinence of the proposed control law.

VIII. CONCLUSION

We consider boundary control of nonlinear ODE/wave PDE cascaded systems with spatially-varying propagation speed. A nonlinear backstepping compensator is designed such that the closed-loop system is globally asymptotically stable and the stability proof is established based on a Lyapunov-like



Fig. 6: Response of wave PDE dynamics under the proposed control

argument. The proposed design is illustrated by the benchmark system controlled via a cable. The generalization of the result to actuator dynamics governed by an arbitrary number of coupled linear hyperbolic PDEs [17], the extension to drilling systems' stabilization problems and to the particular case v(x) = 0 will be considered in our future works.

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